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PREVIOUS YEAR QUESTION

PHYSICAL SCIENCES

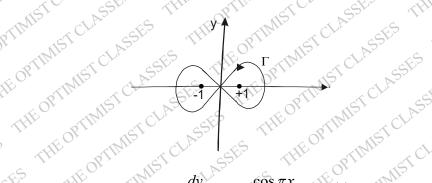
Part - B

- Which of the following cannot be eigenvalues of a real 3×3 matrix
- (b) 1, 1, 1
- (c) $e^{i\theta}$, $e^{-i\theta}$, 1
- (d) i, 1, 0
- (a) 2i, 0, -2i
 22. Let Let $u(x, y) = e^{ax} \cos(by)$ be the real part of a function f(z) = u(x, y) + iv(x, y) of the complex variable z = x + iy, where a, b are real constants and $a \neq 0$. The function f(z) is complex analytic every where in the complex plane if and only if
- (b) $b = \pm a$

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- The integral $\oint \frac{ze^{i\pi z/2}}{}$ -dz along the closed contour Γ shown in the figure is





- The value of Surfice Purific Printing of Surfice Purific Printing (c) -2π (d) $4\pi i$ The function y(x) satisfies the differential equation $x\frac{dy}{dx}$ $2y = \frac{\cos \pi x}{2}$ y(2) is P_{11}

- (d) 1/4
- (c) 1/2 TIMES The random variable $x(-\infty < x < \infty)$ is distributed according to the normal distribution
 - The probability density of the random variable $y = x^2$ is
- $)\frac{1}{2\sqrt{2\pi\sigma^2 v}}e^{\frac{v}{2\sigma^2}},$

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(c)
$$\frac{1}{\sqrt{2\sigma^2}}e^{\frac{y}{2\sigma^2}}, 0 \le y < \infty$$

(d)
$$\frac{1}{\sqrt{2\pi\sigma^2 y}} e^{-\frac{\hat{y}}{\sigma^2}}, 0 \le y < \infty$$

The Hamiltonian for a system described by the generalized coordinate x and generalized momentum p is

$$H = \alpha x^{2} p + \frac{p^{2}}{2(1+2\beta x)} + \frac{1}{2}\omega^{2} x^{2}$$

where α, β and ω are constant. The corresponding Lagrangian is

(a)
$$\frac{1}{2}(\dot{x} - \alpha x^2)(1 + 2\beta x) - \frac{1}{2}\omega^2 x^2$$

(b)
$$\frac{1}{2(1+2\beta x)}\dot{x}^2 - \frac{1}{2}\omega^2 x^2 - \alpha x^2 \dot{x}$$

(c)
$$\frac{1}{2}(\dot{x}^2 - \alpha^2 x)^2 (1 + 2\beta x) - \frac{1}{2}\omega^2 x^2$$

(a)
$$\frac{1}{2}(\dot{x} - \alpha x^2)(1 + 2\beta x) - \frac{1}{2}\omega^2 x^2$$
 (b) $\frac{1}{2(1 + 2\beta x)}\dot{x}^2 - \frac{1}{2}\omega^2 x^2 - \alpha x^2 \dot{x}$ (c) $\frac{1}{2}(\dot{x}^2 - \alpha^2 x)^2(1 + 2\beta x) - \frac{1}{2}\omega^2 x^2$ (d) $\frac{1}{2(1 + 2\beta x)}\dot{x}^2 - \frac{1}{2}\omega^2 x^2 + \alpha x^2 \dot{x}$
An inertial observer sees two events F and F happening at the same location but 6 μ s

- An inertial observer sees two events E_1 and E_2 happening at the same location but $6 \mu s$ apart in time. Another observer moving with a constant velocity v (with respect to the first one) sees the same events to be 9 μs apart. The spatial distance between the events, as measured by the second observer, is approximately
 - $(a) 300 \, m$
- (b) 1000 m
- (c) $2000 \, \text{m}$
- (d) $2700 \, \text{m}$
- A ball weighing 100gm, released from a height of 5 m, bounces perfectly elastically off a plate. The collision time between the ball and the plate is 0.5 s. The average force on the plate is approximately
- A solid vertical rod, of length L, and cross-sectional area A, is made of a material of Young's modulus Y. The rod is loaded with a mass M, and as a result, extends by a small amount ΔL in the equilibrium condition. The mass is then suddenly reduced to M/2. As a result the rod will undergo longitudinal oscillation with an angular frequency
- If the root-mean-squared momentum of a particle in the ground state of a one-dimensional simple harmonic potential is $\,p_0\,$, then its root-mean-squared momentum in the first excited state is

(a)
$$p_0 \sqrt{2}$$

(b)
$$p_0 \sqrt{3}$$

(c)
$$p_0 \sqrt{2/3}$$

(d)
$$p_0 \sqrt{3/2}$$

Consider a potential barrier A of height V_0 and width b, and another potential barrier B of height $2V_0$ and the same width b. The ratio T_A/T_B of tunnelling probabilities T_A and T_B , through barrier A and B respectively, for a particle of energy $V_0 / 100$ is best approximated by

(a)
$$\exp\left[\sqrt{1.99} - \sqrt{0.99}\sqrt{\frac{8mV_0b^2}{\hbar^2}}\right]$$

(b)
$$\exp\left[\sqrt{1.98} - \sqrt{0.98}\sqrt{\frac{8mV_0b^2}{\hbar^2}}\right]$$

(c)
$$\exp\left[\sqrt{2.99} - \sqrt{0.99}\sqrt{\frac{8mV_0b^2}{\hbar^2}}\right]$$

(d)
$$\exp\left[\sqrt{2.98} - \sqrt{0.98}\sqrt{\frac{8mV_0b^2}{\hbar^2}}\right]$$

A constant perturbation H' is applied to a system for time Δt (where $H'\Delta t \ll \hbar$) leading to a

transition from a state with energy E_i to another with energy E_f . If the time of application is doubled, the probability of transition will be

- (a) unchanged
- (c) quadrupled
- (d) halved

- The two vectors $\begin{pmatrix} a \\ 0 \end{pmatrix}$ and $\begin{pmatrix} b \\ c \end{pmatrix}$ are orthogonal if
 - (a) $a = \pm 1, b = \pm 1/\sqrt{2}, c = \pm 1/\sqrt{2}$

(c) $a = \pm 1, b = 0, c = \pm 1$

- (d) $a = \pm 1, b = \pm 1/2, c = 1/2$
- Two long hollow co-axial conducting cylinders of radii R_1 and R_2 ($R_1 < R_2$) are placed in vacuum as shown in the figure below.



The inner cylinder carries a charge $+\lambda$ per unit length and the outer cylinder carries a charge $-\lambda$ per length. The electrostatic energy per unit length of this system is

- (a) $\frac{\lambda^2}{\pi \in_0} \ln \left(\frac{R_2}{R_1} \right)$ (b) $\frac{\lambda^2}{4\pi \in_0} \left(\frac{R_2^2}{R_1^2} \right)$ (c) $\frac{\lambda^2}{4\pi \in_0} \ln \left(\frac{R_2}{R_1} \right)$ (d) $\frac{\lambda^2}{2\pi \in_0} \ln \left(\frac{R_2}{R_1} \right)$
- A set N concentric circular loops of wire, each carrying a steady current I in the same direction, is arranged in a plane. The radius of the first loop is $r_1 = a$ and the radius of the n^{th} loop is given by $r_n = n r_{n-1}$. The magnitude B of the magnetic field at the centre of the circles in the limit $N \to \infty$, is
 - (a) $\frac{\mu_0 I(e^2-1)}{4\pi a}$ (b) $\frac{\mu_0 I(e-1)}{\pi a}$ (c) $\frac{\mu_0 I(e^2-1)}{8a}$ (d) $\frac{\mu_0 I(e-1)}{2a}$

- An electromagnetic wave (of wavelength λ_0 in free space) travels through an absorbing medium with dielectric permittivity given by $\varepsilon = \varepsilon_R + i\varepsilon_I$, where $\frac{\varepsilon_I}{\varepsilon_R} = \sqrt{3}$. If the skin depth is $\frac{\Lambda_0}{4\pi}$, the ratio of the amplitude of electric field E to that of the magnetic field B, in the medium (in ohms) is
 - (a) $120 \, \pi$
- (c) $30\sqrt{2} \pi$
- (d) 30 π
- The vector potential $\vec{A} = ke^{-at}r\hat{r}$, (where a and k are constants) corresponding to an electromagnetic field is changed to $\vec{A}' = -ke^{-at}r\hat{r}$. This will be a gauge transformation if the corresponding change $\phi' - \phi$ in the scalar potential is
 - (a) akr^2e^{-at}

- A thermodynamic function G(T,P,N) = U TS + PV is given in terms of the internal energy U, temperature T, entropy S, pressure P, volume V and the number of particles N. Which of the following relations is

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true? (In the following μ is the chemical potential).

(a)
$$S = -\frac{\partial G}{\partial T}\Big|_{N}$$

(b)
$$S = \frac{\partial G}{\partial T}\Big|_{M = 0}$$

(c)
$$V = -\frac{\partial G}{\partial P}\Big|_{N,T}$$

(d)
$$\mu = -\frac{\partial G}{\partial N}\Big|_{P}$$

A box, separated by a movable wall, has two compartments filled by a monoatomic gas of $\frac{C_p}{C} = \gamma$

Initially the volumes of the two compartments are equal, but the pressure are $3p_0$ and p_0 respectively. When the wall is allowed to move, the final pressure in the two compartments becomes equal. The final pressure is

(a)
$$\left(\frac{2}{3}\right)^{\gamma} p_0$$

(b)
$$3\left(\frac{2}{3}\right)^{x} p_0$$

(c)
$$\frac{1}{2} \left(1 + 3^{1/x} \right)^{x} p$$

(d)
$$\left(\frac{3^{1/\gamma}}{1+3^{1/\gamma}}\right)^{\gamma} p_0$$

- A gas of photons inside a cavity of volume V is in equilibrium at temperature T. If the temperature of the cavity is changed to 2T, the radiation pressure will change by a factor of (a) 2

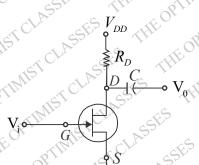
- 41. In a thermodynamic system in equilibrium, each molecule can exist in three possible states with probabilities $\frac{1}{2}$, $\frac{1}{3}$ and $\frac{1}{6}$ respectively. The entropy per molecule is

(a)
$$k_B \ln 3$$

(b)
$$\frac{1}{2}k_B \ln 2 + \frac{2}{3}k_B \ln 3$$

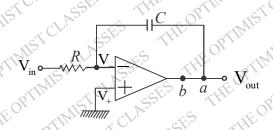
(a)
$$k_B \ln 3$$
 (b) $\frac{1}{2}k_B \ln 2 + \frac{2}{3}k_B \ln 3$ (c) $\frac{2}{3}k_B \ln 2 + \frac{1}{2}k_B \ln 3$ (d) $\frac{1}{2}k_B \ln 2 + \frac{1}{6}k_B \ln 3$ In the *n*-channel JFET shown in figure below,

(d)
$$\frac{1}{2}k_B \ln 2 + \frac{1}{6}k_B \ln 3$$



If the drain D-source S saturation current I_{DSS} is 10 mA and the pinch-off voltage V_p is -8V, then the voltage across points D and S is (a) 11.125V (b) 10.375V (c) 5.75V (d) 4.75V

The gain of the circuit given below is =



The modification in the circuit required to introduce a dc feedback is to add a resistor

- (a) between a and b
- (b) between positive terminal of the op-amp and ground

(c) in series with C

- (d) parallel to C
- A 2×4 decoder with an enable input can function as a
 - (a) 4×1 multiplexer

b) 1×4 demultiplexer

(c) 4×2 encoder

- (d) 4×2 priority encoder
- The experimentally measured values of the variables x and y are 2.00 ± 0.05 and 3.00 ± 0.02 , respectively.

 What is the error in the calculated value of $x = 2.00 \pm 0.05$. What is the error in the calculated value of z = 3y - 2x from the measurements?
- (c) 0.03 Part - C THE OPTIM

- The Green's function satisfying $\frac{d^2}{dx^2}g(x,x_0) = \delta(x-x_0)$, with the boundary conditions is
- $\begin{cases} \frac{1}{2L} (x_0 L)(x + L), -L \le x < x_0 \end{cases}$
- $\begin{cases} \frac{1}{2L}(x_0 + L)(x + L), -L \le x < x_0 \\ \frac{1}{2T}(x_0 L) \end{cases}$ $\frac{1}{2L}(x_0 - L)(x - L)$
- $(d) \frac{1}{2L}(x-L)(x+L)$
- 7. Let σ_x, σ_y and σ_z be the pauli matrices and

(a)
$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta & 0 \\ \sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

$$\begin{aligned}
\text{(b)} \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} &= \begin{pmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} \\
\begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} & 0 \\ 0 & x \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

(c)
$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \cos\frac{\theta}{2} & \sin\frac{\theta}{2} & 0 \\ -\sin\frac{\theta}{2} & \cos\frac{\theta}{2} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

$$(d) \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \cos\frac{\theta}{2} & -\sin\frac{\theta}{2} & 0 \\ \sin\frac{\theta}{2} & \cos\frac{\theta}{2} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

8. The interval [0, 1] is divided into 2n parts of equal length to calculate the integral $\int_{0}^{1} e^{i2\pi x} dx$ using Simpson's

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$\frac{1}{2}$ -rule.	What is the	minimum	value of n	for the res	ult to be ex	act?
3 Miles	CLAD.	~S	OX	MSI	at AS	, 6

- 49. Which of the following sets of 3×3 matrices (in which a and b are real numbers) a group under matrix multiplication?
 - (a) $\left\{ \begin{pmatrix} 1 & 0 & a \\ 0 & 1 & 0 \\ b & 0 & 1 \end{pmatrix} : a, b \in \mathbb{R} \right\}$
- (b) $\left\{ \begin{pmatrix} 1 & a & 0 \\ 0 & 1 & b \\ 0 & 0 & 1 \end{pmatrix} : a, b \in \mathbb{R} \right\}$

(c) $\left\{ \begin{pmatrix} 1 & 0 & a \\ 0 & 1 & b \\ 0 & 0 & 1 \end{pmatrix} : a, b \in \mathbb{R} \right\}$

- (d) $\left\{ \begin{pmatrix} 1 & a & 0 \\ b & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} : a, b \in \mathbb{R} \right\}$
- The Lagrangian of a free relativistic particle (in one-dimension) of mass m is given by $L = -m\sqrt{1-\dot{x}^2}$, where $\dot{x} = dx/dt$. If such a particle is acted upon by a constant force in the direction of its motion, the phase space trajectories obtained from the corresponding Hamiltonian are

 (a) ellipses

 (b) cycloids

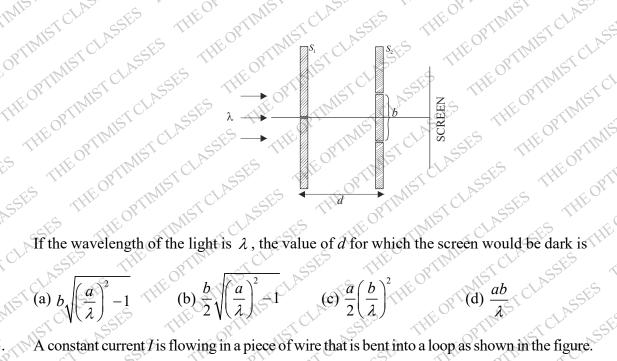
 (c) hyperbolas

 (d) parabolas
- 51. A Hamiltonian system is described by the canonical coordinate q and canonical momentum p. A new coordinate Q is defined as $Q(t) = q(t+\tau) + p(t+\tau)$, where t is the time and τ is a constant, that is, the new coordinate is a combination of the old coordinate and momentum at a shifted time. The new canonical momentum p(t) can be expressed as
 - (a) $p(t+\tau)-q(t+\tau)$

(b) $p(t+\tau)-q(t-\tau)$

(c) $\frac{1}{2} [p(t-\tau)-q(t+\tau)]$

- (d) $\frac{1}{2} [p(t+\tau) q(t+\tau)]$
- 52. The energy of a one-dimensional system, governmend by the Lagrangian $L = \frac{1}{2}m\dot{x}^2 \frac{1}{2}kx^{2n}$, where k and n are two positive constants, is E_0 . The time period of oscillation τ satisfies
 - (a) $\tau \propto k^{-1/n}$
- (b) $\tau \propto k^{-1/2n} E_0^{\frac{1-n}{2n}}$
- (c) $\tau \propto k^{-1/2n} E_0^{\frac{n-2}{2n}}$
- (d) $\tau \propto k^{-1/n} E_0^{\frac{1+n}{2n}}$
- 53. An electron is decelerated at a constant rate starting from an initial velocity u (where $u \ll c$) to u/2
 - during which it travels a distance s. The amount of energy lost to radiation is
 - (a) $\frac{\mu_0 e^2 u^2}{3\pi mc^2 s}$
- (b) $\frac{\mu_0 e^2 u^2}{6\pi mc^2 s}$
- (c) $\frac{\mu_0 e^2 u}{8\pi mcs}$
- (d) $\frac{\mu_0 e^2 u}{16\pi mcs}$
- The figure below describes the arrangement of slits and screens in a Young's double slit experiment. The width of the slit in S_1 is a and the slits in S_2 are of negligible width.



(a)
$$b\sqrt{\left(\frac{a}{\lambda}\right)^2-1}$$

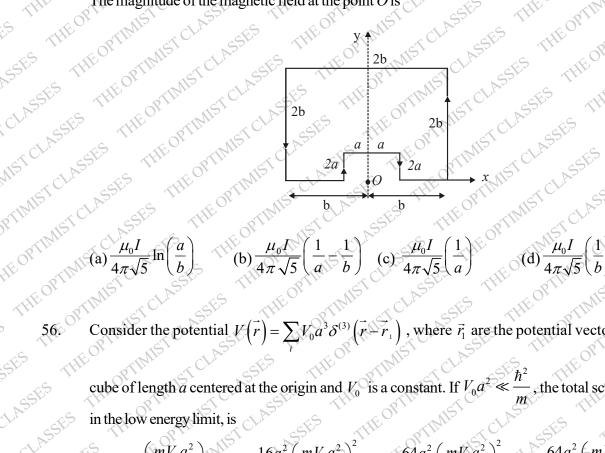
(b)
$$\frac{b}{2}\sqrt{\left(\frac{a}{\lambda}\right)^2-1}$$

(c)
$$\frac{a}{2} \left(\frac{b}{\lambda} \right)^2$$

(d)
$$\frac{ab}{\lambda}$$

55. A constant current *I* is flowing in a piece of wire that is bent into a loop as shown in the figure.

The magnitude of the magnetic field at the part of the magnitude of the magnetic field at the part of t The magnitude of the magnetic field at the point O is



(a)
$$\frac{\mu_0 I}{4\pi\sqrt{5}} \ln\left(\frac{a}{b}\right)$$

$$\text{(b)} \frac{\mu_0 I}{4\pi\sqrt{5}} \left(\frac{1}{a} - \frac{1}{b} \right)$$

(c)
$$\frac{\mu_0 I}{4\pi\sqrt{5}} \left(\frac{1}{a}\right)$$

$$(d) \frac{\mu_0 I}{4\pi\sqrt{5}} \left(\frac{1}{b}\right)$$

(a) $\frac{\mu_0 I}{4\pi\sqrt{5}} \ln\left(\frac{a}{b}\right)$ (b) $\frac{\mu_0 I}{4\pi\sqrt{5}} \left(\frac{1}{a} = \frac{1}{b}\right)$ (c) $\frac{\mu_0 I}{4\pi\sqrt{5}} \left(\frac{1}{a}\right)$ (d) $\frac{\mu_0 I}{4\pi\sqrt{5}} \left(\frac{1}{b}\right)$ Consider the potential $V(\vec{r}) = \sum_i V_0 a^3 \delta^{(3)} \left(\vec{r} - \vec{r}_1\right)$, where \vec{r}_1 are the potential vectors of the vertices of a

cube of length a centered at the origin and V_0 is a constant. If $V_0 a^2 \ll \frac{\hbar^2}{m}$, the total scattering cross-section, in the low energy limit, is $(a) \ 16a^2 \left(\frac{mV_0 a^2}{\hbar^2}\right) \qquad (b) \ \frac{16a^2}{\pi^2} \left(\frac{mV_0 a^2}{\hbar^2}\right)^2 \qquad (c) \ \frac{64a^2}{\pi} \left(\frac{mV_0 a^2}{\hbar^2}\right)^2 \qquad (d) \ \frac{64a^2}{\pi^2} \left(\frac{mV_0 a^2}{\hbar^2}\right)$

(a)
$$16a^2 \left(\frac{mV_0a^2}{\hbar^2}\right)$$

(b)
$$\frac{16a^2}{\pi^2} \left(\frac{mV_0 a^2}{\hbar^2} \right)$$

(c)
$$\frac{64a^2}{\pi} \left(\frac{mV_0 a^2}{\hbar^2} \right)$$

(d)
$$\frac{64a^2}{\pi^2} \left(\frac{mV_0 a^2}{\hbar^2} \right)$$

57. The Coulomb potential $V(r) = -e^2/r$ of a hydrogen atom is perturbed by adding $H' = bx^2$ (where b is a constant) to the Hamiltonian. The first order correction to the ground state. is a constant) to the Hamiltonian. The first order correction to the ground state energy is (The ground state wavefunction is $w = \frac{1}{2} e^{-r/ac}$.

	1
	01 2
(8)	$2ba_{0}^{2}$
(ha)	$20u_0$

(b)
$$2ba_0^2$$

(c)
$$ba_0^2 / 2$$

(d)
$$\sqrt{2}ba_0^2$$

(a)
$$2ba_0^2$$
 (b) $2ba_0^2$ (c) $ba_0^2/2$ (d) $\sqrt{2}$

58. Using the trial function $\psi(x) = \begin{cases} A(a^2 - x^2); & -a < x < a \\ 0 & ; \text{ otherwise} \end{cases}$ (e) $\frac{1}{2}\hbar\omega$ (d) $\sqrt{\frac{5}{7}}$

(b)
$$\sqrt{\frac{5}{14}}\hbar\omega$$

(c)
$$\frac{1}{2}\hbar\omega$$

(d)
$$\sqrt{\frac{5}{7}}\hbar\omega$$

m, consider In the usual notation $|nlm\rangle$ for the states of a hydrogen like atom, consider the spontaneous transitions $|210\rangle \rightarrow |100\rangle$ and $|310\rangle \rightarrow |100\rangle$. If t_1 and t_2 are the lifetimes of the first and the second decaying states respectively, then the ratio t_1 / t_2 is proportional to

(a)
$$\left(\frac{32}{27}\right)^3$$

(b)
$$\left(\frac{27}{32}\right)^3$$

(c)
$$\left(\frac{2}{3}\right)^3$$

(d)
$$\left(\frac{3}{2}\right)^3$$

A random variable n obeys poisson statistics. The probability of finding n = 0 is 10 %. The expectation value of

The single particle energy levels of a non-interacitng three-dimensional isotropic system, labelled by momen tum k, are proportional to k^3 . The ratio P/ε of the average pressure \vec{P} to the energy density ε at a fixed

(a) 1/3

(b)
$$2/3$$

The Hamiltonian for three Ising spins S_0, S_1 and S_2 taking values ± 1 is $H = -j S_0 (S_1 + S_2)$. If the system is in equilibrium at temperature T, the average energy of the system, in terms of $\beta = (k_B T)^{-1}$

(a)
$$-\frac{1+\cosh(2\beta j)}{2\beta\sinh(2\beta j)}$$

(b)
$$-2j\left[1+\cosh\left(2\beta j\right)\right]$$

(c)
$$-2/\beta$$

(d)
$$-2j \frac{\sinh(2\beta j)}{1+\cosh(2\beta j)}$$

Let I_0 be the saturation current, η the ideality factor and v_F and v_R the forward and reverse poten

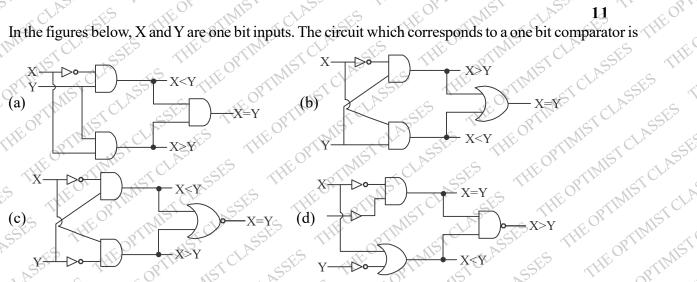
tials, respectively, for a diode. The ratio R_R / R_F of its reverse and forward resistances R_R and R_F (b) $\frac{v_E}{v_R} \exp\left(\frac{qv_E}{\eta k_B T}\right)$ respectively, varies as (In the following k_B is the Boltzmann constant, T is the absolute temperature and q is the

(a)
$$\frac{v_R}{v_F} \exp\left(\frac{qv_F}{\eta k_B T}\right)$$

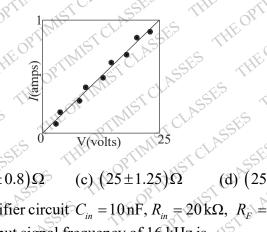
(b)
$$\frac{v_F}{v_R} \exp\left(\frac{qv_F}{\eta k_B T}\right)$$

(c)
$$\frac{v_R}{v_F} \exp\left(-\frac{qv_F}{\eta k_B T}\right)$$

(d)
$$\frac{v_F}{v_R} \exp \left(-\frac{qv_F}{\eta k_B T}\right)$$

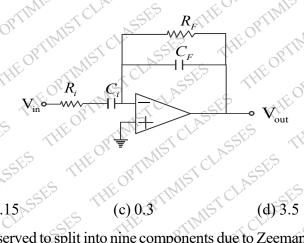


Both the data points and a linear fit to the current vs voltage of a resistor are shown in the graph below. If the error in the slope is $1.255 \times 10^{-3} \, \Omega^{-1}$, then the value of resistance estimated from the graph is



- (a) $(0.04 \pm 0.8)\Omega$ (b) $(25.0 \pm 0.8)\Omega$

(d) $(25\pm0.0125)\Omega$ Continued Ω , $R_E = 200 \text{ k}\Omega$ In the following operational amplifier circuit $C_{in}=10\,\mathrm{nF}$, $R_{in}=20\,\mathrm{k}\Omega$, $R_F=200\,\mathrm{k}\Omega$ and $C_F=100\,\mathrm{pF}$. The magnitude of the gain at a input signal frequency of 16 kHz is The magnitude of the gain at a input signal frequency of 16 kHz is



- (d) 3.5

67. ASSE Ar ME(b) 0.15 5515 An atomic spectral line is observed to split into nine components due to Zeeman shift. If the upper state of the atom is ³D, then the lower state will be the atom is 3D_2 then the lower state will be

- (d) 3 P.

(a) 3 F₂ If the coefficient of stimulated emission for a particular transition is $2.1 \times 10^{19} \text{ m}^3 \text{W}^$ emitted photon is at wavelength 3000 Å, then the lifetime of the excited sate is approximately.

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OPTIMIS	(a) 20 ns	(b) 40 ns	(c) 80 ns	(d) 100 ns	12 THE OP I
5 THE 69. OPT	If the bindings energies respectively, then the	s of the electron in the	e K and L shells of silv	ver atom are 25.4 keV and be approximately	3.34 keV,
3SES THE	(a) 22 keV	(b) 9.3 keV	(c) 10.5 keV	(d) 18.7 keV	SI CLASO SES II
	respectviely. For simp by broadband radiatio	licity take the dielec n, an electron initial	tric constant of the nly in the valence ban	miconductor are 1.875 eV naterial to the unity. Whe d at $k = 0$ makes a transit ction band, in terms of th	n it is excited ion to the
181 CLASE	k_{max} at the edge of the	Brillouin zone, afte	er the transition is clo	osest to	THE OF I. TIMISTE
TIMIS IST CLA			(c) $k_{\text{max}}/1000$	TIME (d) OCLASSIS	THE OP I. PETIMISTS
E TIME	while the electron den	sity in silver is appro	oximately 10% of the	the electrical conductivit e electron density in copp	er. In Drude's
THEOR	model, the approxima	te ratio $ au_{Cu}$ / $ au_{Ag}$ of t	he mean collision tin	me in copper (au_{cu}) to the	mean collision
THEOR	time in silver $\left(au_{{\scriptscriptstyle A} {\scriptscriptstyle g}} \right)$ is	S THE OP .	IST CLASS SES	me in copper (τ_{eu}) to the	CLASSI THE
JES THE	(a) 0.44	(b) 1.50	(c) 0.33	SES (d) 0.66 Primisi Cr	CLAN
SSES 72.	The charge distribution	n inside a material o	f conductivity σ and	d permittivity ε at initia	I time $t = 0$ is
LALSSES	$\rho(r,0) = \rho_0$, constant	t. At subsequent time	es $\rho(r,t)$ is given b	YAS SES THEORY	TIMIST CLASS SE
MSTCLASSES	(a) $\rho_0 \exp\left(-\frac{\sigma t}{\varepsilon}\right)$ (c) $\rho_0 \left(\frac{\rho_0}{\sigma t}\right)$	IST CLASSES	(b) $\frac{1}{2}\rho_0 \left[1 + \exp \left(\frac{1}{2}\rho_0\right)\right]$	d permittivity ε at initially $0 \left(\frac{\sigma t}{\varepsilon} \right)$	Itime t=0 is PUTINIST CLASSIS FILE OPTIMIST CLASSIS
OPTIMIST CLASS OPTIMIST CLASS OPTIMIST CLASS OPTIMIST CLASS	(c) $\frac{\rho_0}{\left[1 - \exp\left(\frac{\sigma t}{\varepsilon}\right)\right]}$	FOPTIMIST CLASSES	(d) $\rho_0 \cosh \frac{\sigma \iota}{c}$	released in the reaction is $(d) \frac{23Q}{58}$ we exchange of pions is 1	THE STIP
73. NE	If in a spontaneous α	-decay of $\frac{232}{92}$ U at r	est, the total energy	released in the reaction is	Q, then the
CIF OF I	energy carried by the	lpha -particle is	STO CLASSIT	THE STOPPING OF STOP	ASSES STILL
S THE OPT	(a) $\frac{5/Q}{58}$	(b) $\frac{Q}{57}$	$\frac{Q}{58}$ LASSIT	(d) $\frac{23Q}{58}$ The exchange of pions is 1.	then the THE OF THE OF THE OF THE OF
74. THE	The range of the nucle mass of the pion is 14	ear force between tw 40 MeV/c ² and the m	o nucleons due to the nass of the rho-meso	the exchange of pions is 1. In is 770 MeV/c^2 , then the	40 fm. If the erange of the
CLASSIS	force due to exchange (a) 1.40 fm	of rho mesons is (b) 7.70 fm	(c) 0.25 fm	(d) 0.18 fm	OPTIMIST CLASSIC
75.ASS	A baryon X decays by s	strong interaction as	$X \to \Sigma^+ + \pi \bar{} + \pi^0$, W	where Σ^+ is a member of the	ne isotriplet
AFTIMIST CLASSINGST CL	$(\Sigma^+, \Sigma^0, \Sigma^-)$. The thire (a) 0	d component I_3 of the (b) $1/2$	e isospin of X is	the exchange of pions is 1.5 m is 770 MeV/c ² , then the (d) 0.18 fm where Σ^+ is a member of the (d) 3/2	THE OPTIMEST OF THE OPTIMEST
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